

1. **Solution:**

The Lagrangian for the whole system is the sum of the Lagrangians  $L = L_1 + L_2$ : mass point 1 (mass  $m_1$ ) and mass point 2 (mass  $m_2$ ).

**particle 1**

Particle 1 (the attachment point) has the location

$$\mathbf{r}_1 = x\mathbf{i}$$

and thus the kinetic energy

$$T_1 = \frac{1}{2}m_1\dot{x}^2.$$

We can choose that the particle 1 has potential  $V_1 = 0$ .

**particle 2**

Particle 2 (the pendulum) has the location

$$\mathbf{r}_2 = (x + l \sin \phi)\mathbf{i} - l \cos \phi\mathbf{j}$$

and the velocity

$$\mathbf{v}_2 = \dot{\mathbf{r}}_2 = (\dot{x} + l \cos \phi \dot{\phi})\mathbf{i} + l \sin \phi \dot{\phi}\mathbf{j}.$$

Thus the kinetic energy has the form

$$\begin{aligned} T_2 &= \frac{1}{2}m_2v_2^2 \\ &= \frac{1}{2}m_2\mathbf{v}_2 \cdot \mathbf{v}_2 \\ &= \frac{1}{2}m_2[(\dot{x} + l \cos \phi \dot{\phi})^2 + l^2 \sin^2 \phi \dot{\phi}^2] \\ &= \frac{1}{2}m_2[\dot{x}^2 + 2l\dot{x}\dot{\phi} \cos \phi + l^2 \cos^2 \phi \dot{\phi}^2 + l^2 \sin^2 \phi \dot{\phi}^2] \\ &= \frac{1}{2}m_2[\dot{x}^2 + 2l\dot{x}\dot{\phi} \cos \phi + l^2 \dot{\phi}^2]. \end{aligned}$$

When we calculate the potential energy we have to remember our earlier choice. Then the potential is

$$V_2 = -m_2gl \cos \phi.$$

Now the Lagrangian for the whole system is

$$\begin{aligned} L &= L_1 + L_2 \\ &= T_1 + T_2 - (V_1 + V_2) \\ &= \frac{1}{2}(m_1 + m_2)\dot{x}^2 + m_2l\dot{x}\dot{\phi} \cos \phi + \frac{1}{2}m_2l^2\dot{\phi}^2 + m_2gl \cos \phi. \end{aligned}$$

## 2. Solution:

The particle has the position

$$\mathbf{r}_1 = [a \cos \omega t + l \sin \phi] \mathbf{i} + [a \sin \omega t - l \cos \phi] \mathbf{j}$$

and thus the velocity is

$$\mathbf{v}_1 = [-a\omega \sin \omega t + l\dot{\phi} \cos \phi] \mathbf{i} + [a\omega \cos \omega t + l\dot{\phi} \sin \phi] \mathbf{j}.$$

meaning that the kinetic energy is

$$\begin{aligned} T &= \frac{1}{2} m \mathbf{v} \cdot \mathbf{v} \\ &= \frac{1}{2} m [(-a\omega \sin \omega t + l\dot{\phi} \cos \phi)^2 + (a\omega \cos \omega t + l\dot{\phi} \sin \phi)^2] \\ &= \frac{1}{2} m [a^2 \omega^2 \sin^2 \omega t - 2al\omega\dot{\phi} \cos \phi \sin \omega t + l^2 \dot{\phi}^2 \cos^2 \phi \\ &\quad + a^2 \omega^2 \cos^2 \omega t + 2al\omega\dot{\phi} \sin \phi \cos \omega t + l^2 \dot{\phi}^2 \sin^2 \phi] \\ &= \frac{1}{2} m [a^2 \omega^2 \underbrace{(\sin^2 \omega t + \cos^2 \omega t)}_{=1} + 2al\omega\dot{\phi} \underbrace{(\sin \phi \cos \omega t - \cos \phi \sin \omega t)}_{=\sin(\phi-\omega t)} + l^2 \dot{\phi}^2 \underbrace{(\sin^2 \omega t + \cos^2 \omega t)}_{=1}] \\ &= \frac{1}{2} m a^2 \omega^2 + m l a \omega \dot{\phi} \sin(\phi - \omega t) + \frac{1}{2} m l^2 \dot{\phi}^2 \end{aligned}$$

The potential energy is

$$\begin{aligned} V &= mgy, \quad y = a \sin \omega t - l \cos \phi \\ &= mg(a \sin \omega t - l \cos \phi) \\ &= mga \sin \omega t - mgl \cos \phi. \end{aligned}$$

Thus the Lagrangian is

$$\begin{aligned} L &= T - V \\ &= \frac{1}{2} m l^2 \dot{\phi}^2 + m l a \omega \dot{\phi} \sin(\phi - \omega t) + mgl \cos \phi - mga \sin \omega t + \frac{1}{2} m a^2 \omega^2. \end{aligned}$$

Let's denote

$$L_0 = -mga \sin \omega t + \frac{1}{2} m a^2 \omega^2.$$

Now we see that

$$\frac{\partial L_0}{\partial \dot{q}} = \frac{\partial L_0}{\partial q} = 0,$$

where in our case  $q = \phi$ . So  $L_0$  does not contribute to the equation of motion.

## 3. Solution:

Before proving anything, let's recall some useful results:

$$\nabla \times \nabla \phi \equiv 0$$

$$\nabla \cdot (\nabla \times \mathbf{A}) \equiv 0.$$

These results hold for an arbitrary scalar field  $\phi$  and vector field  $\mathbf{A}$ . Of course, we assume that the needed derivatives exist (this is usually the case in physics). If you do not believe, you can prove the results by simple calculations.

a)

Now we have

$$\mathbf{E} = -\nabla\phi - \partial_t\mathbf{A}$$

$$\mathbf{B} = \nabla \times \mathbf{A}.$$

Thus

$$\nabla \cdot \mathbf{B} = \nabla \cdot \nabla \times \mathbf{A} = 0$$

and

$$\begin{aligned} \nabla \times \mathbf{E} &= \nabla \times (-\nabla\phi - \partial_t\mathbf{A}) \\ &= -\underbrace{\nabla \times \nabla\phi}_{=0} - \underbrace{\nabla \times \partial_t\mathbf{A}}_{=\partial_t\nabla \times} \\ &= -\partial_t \underbrace{\nabla \times \mathbf{A}}_{=\mathbf{B}} \\ &= -\partial_t\mathbf{B}. \end{aligned}$$

So the fields  $\mathbf{E}$  and  $\mathbf{B}$  defined by the scalar field  $\phi$  and the vector field  $\mathbf{A}$  produce two of the Maxwell equations

$$\nabla \cdot \mathbf{B} = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial\mathbf{B}}{\partial t}.$$

b)

Now we make a gauge transformation to the potentials i. e.

$$\phi' = \phi - \partial_t\chi$$

$$\mathbf{A}' = \mathbf{A} + \nabla\chi$$

that defines new fields  $\mathbf{E}' = -\nabla\phi' - \partial_t\mathbf{A}'$  and  $\mathbf{B}' = \nabla \times \mathbf{A}'$ . By simple calculations we get

$$\begin{aligned} \mathbf{E}' &= -\nabla\phi' - \partial_t\mathbf{A}' \\ &= -\nabla(\phi - \partial_t\chi) - \partial_t(\mathbf{A} + \nabla\chi) \\ &= -\nabla\phi + \partial_t\nabla\chi - \partial_t\mathbf{A} - \partial_t\nabla\chi, \quad (\partial_t\nabla = \nabla\partial_t) \\ &= -\nabla\phi - \partial_t\mathbf{A} \\ &= \mathbf{E} \end{aligned}$$

and

$$\begin{aligned}\mathbf{B}' &= \nabla \times \mathbf{A}' \\ &= \nabla \times (\mathbf{A} + \nabla\chi) \\ &= \underbrace{\nabla \times \mathbf{A}}_{=\mathbf{B}} + \underbrace{\nabla \times \nabla\chi}_{=0} \\ &= \mathbf{B}.\end{aligned}$$

So we see that the gauge transformation does not change the fields. We call that the fields are *gauge invariants*.

4. **Solution:**

We have a charged particle (the mass  $m$  and the charge  $q$ ) in a magnetic field  $\mathbf{B} = B\mathbf{k}$ . Let's choose a potential as  $\mathbf{A} = -By\mathbf{i}$ . Now

$$\nabla \times \mathbf{A} = \nabla \times (-By)\mathbf{i} = B\mathbf{k} = \mathbf{B}.$$

So our potential gives the correct magnetic field.

The velocity of the particle is  $\mathbf{v} = \dot{x}\mathbf{i} + \dot{y}\mathbf{j} + \dot{z}\mathbf{k}$ . The potential energy using the result from lectures is

$$\begin{aligned}V &= q(\phi - \mathbf{v} \cdot \mathbf{A}) \\ &= -q(\dot{x}\mathbf{i} + \dot{y}\mathbf{j} + \dot{z}\mathbf{k}) \cdot (-By)\mathbf{i} \\ &= qB\dot{x}y\end{aligned}$$

Note that one can always choose the scalar potential to be zero  $\phi \equiv 0$  (because of the result of previous question 3b). Now the kinetic energy

$$T = \frac{1}{2}mv^2 = \frac{1}{2}m(\mathbf{v} \cdot \mathbf{v}) = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2).$$

Thus the Lagrangian of the particle is

$$L = T - V = \frac{1}{2}m(\dot{x}^2 + \dot{y}^2 + \dot{z}^2) - qB\dot{x}y.$$

The equation of motion for the particle is given by the Lagrange equation:

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{q}} \right) - \frac{\partial L}{\partial q} = 0.$$

**x-direction**

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x} = 0 \Leftrightarrow m\ddot{x} - qB\dot{y} = 0$$

**y-direction**

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{y}} \right) - \frac{\partial L}{\partial y} = 0 \Leftrightarrow m\ddot{y} + qB\dot{x} = 0$$

## z-direction

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{z}} \right) - \frac{\partial L}{\partial z} = 0 \Leftrightarrow m\ddot{z} = 0$$

The equation of motion in z-direction is easy to solve:

$$\ddot{z} = 0 \Rightarrow z(t) = z(0) + \dot{z}(0)t.$$

This result means that the particle in the z-direction is moving with a constant velocity. In the xy-plane we have to solve a set of equations:

$$\begin{aligned}\ddot{x} &= \frac{qB}{m}\dot{y} \Rightarrow \ddot{x} = -\omega\dot{y} \\ \ddot{y} &= -\frac{qB}{m}\dot{x} \Rightarrow \ddot{y} = \omega\dot{x}\end{aligned}$$

where we use a notation

$$\omega = -\frac{qB}{m}.$$

The equations can be solved by many ways. One method is

$$\begin{aligned}\ddot{x} &= -\omega\dot{y} \\ \Leftrightarrow \\ \frac{d}{dt}\dot{x} &= -\omega\dot{y}, \quad \ddot{y} = \omega\dot{x} \\ \Leftrightarrow \\ \ddot{x} &= -\omega^2\dot{x}, \quad \text{notation } v_x = \dot{x} \\ \Leftrightarrow \\ \ddot{v}_x + \omega^2 v_x &= 0, \quad \text{harmonic oscillator} \\ \Leftrightarrow \\ v_x &= A \sin(\omega t + \phi_0) \\ x &= \int v_x dt = -\frac{A}{\omega} \cos(\omega t + \phi_0) + x_0\end{aligned}$$

and thus

$$\begin{aligned}\ddot{x} &= -\omega\dot{y} \\ \Leftrightarrow \\ \dot{y} &= -\frac{1}{\omega}\ddot{x} = -A \cos(\omega t + \phi_0) \\ \Leftrightarrow \\ y &= \int \dot{y} dt = -\frac{A}{\omega} \sin(\omega t + \phi_0) + y_0.\end{aligned}$$

Now let's denote

$$r_0 \equiv -\frac{A}{\omega}.$$

Thus our solution is

$$x = r_0 \cos(\omega t + \phi_0) + x_0$$

$$y = r_0 \sin(\omega t + \phi_0) + y_0.$$

where these six free parameters  $r_0$ ,  $\phi_0$ ,  $x_0$ ,  $y_0$ ,  $z(0)$  and  $\dot{z}(0)$  are determined by the initial values.