

Answering correctly to 4 questions gives full points; however, answering to all 5 contributes to the total point score.

1. Let us consider (non-relativistic) harmonic oscillator with

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2\hat{x}^2,$$

and define operators

$$\hat{a} = \sqrt{\frac{m\omega}{2}}\hat{x} + \frac{i}{\sqrt{2m\omega}}\hat{p} \quad \hat{a}^\dagger = \sqrt{\frac{m\omega}{2}}\hat{x} - \frac{i}{\sqrt{2m\omega}}\hat{p}.$$

Using

$$[\hat{x}, \hat{p}] = i, \quad [\hat{x}, \hat{x}] = [\hat{p}, \hat{p}] = 0$$

show that

- a) $[\hat{a}, \hat{a}] = [\hat{a}^\dagger, \hat{a}^\dagger] = 0$
 - b) $[\hat{a}, \hat{a}^\dagger] = 1,$
 - c) $\hat{H} = \omega(\hat{a}^\dagger\hat{a} + \frac{1}{2})$
2. Let $c = (c_1 \ c_2)^T$, $c^\dagger = (c_1^* \ c_2^*)$ be 2-component Grassmann vectors, and M a 2×2 complex matrix. Calculate the Grassmann integral

$$\int dc_1^* dc_1 dc_2^* dc_2 \exp[-c^\dagger M c]$$

3. Show that the solution of the non-homogeneous equation ($J(x)$: source term)

$$(\partial^\mu\partial_\mu + m^2)\psi(x) = J(x)$$

can be written as

$$\psi(x) = \psi_0(x) + i \int d^4y G(x, y) J(y)$$

where ψ_0 is a solution for the homogeneous equation and G is a Green function. Convince yourself that if G is the retarded propagator, $J(y)$ can affect $\psi(x)$ only if y is in the past light cone of point x (i.e. $y^0 \leq x^0$, $(x - y)^2 \geq 0$).

4. Let us consider the “scalar quantum electrodynamics”, a theory with a complex scalar field $\phi(x)$ and an abelian gauge field $A_\mu(x)$. The action is

$$S = \int d^4x \left[-\frac{1}{4}F_{\mu\nu}(x)F^{\mu\nu}(x) + (D_\mu(x)\phi(x))^* D^\mu(x)\phi(x) - m^2\phi^*(x)\phi(x) \right], \quad (1)$$

where

$$D_\mu(x) = \partial_\mu - ieA_\mu(x), \quad F_{\mu\nu}(x) = \partial_\mu A_\nu(x) - \partial_\nu A_\mu(x),$$

and e is the coupling constant.

Show that the action is invariant under the gauge transformation ($\theta(x)$ a real number)

$$\phi(x) \rightarrow e^{i\theta(x)}\phi(x), \quad \phi^*(x) \rightarrow e^{-i\theta(x)}\phi^*(x), \quad A_\mu(x) \rightarrow A_\mu(x) - \frac{1}{e}\partial_\mu\theta(x).$$

5. In scalar quantum electrodynamics (eq.(1)),

- a) write down the Feynman propagator $G(p)$ for the scalar field ϕ .
- b) For the free gauge field, show that we can write

$$\int d^4x -\frac{1}{4}F_{\mu\nu}(x)F^{\mu\nu}(x) = \int d^4x d^4y \frac{1}{2}A_\mu(x)[g^{\mu\nu}\partial_\alpha\partial^\alpha - \partial^\mu\partial^\nu]\delta^4(x-y)A_\nu(y).$$

- c) Using the Lorenz gauge $\partial_\nu A^\nu = 0$, argue that the gauge field Feynman propagator is

$$D_{\mu\nu}(q) = \frac{-ig_{\mu\nu}}{q^2 + i\epsilon}.$$

(Note that in this case you need not to use the gauge parameter ξ , but can impose the gauge condition directly. Ignore ghosts.)